



## Molecular Crystals and Liquid Crystals Science and Technology. Section A. Molecular Crystals and Liquid Crystals

Publication details, including instructions for authors and subscription information:

<http://www.tandfonline.com/loi/gmcl19>

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Version of record first published: 04 Oct 2006

To cite this article: L. Duclaux, F. Beguin, K. Andersen & B. Ouladdiaf (1998): Neutron Diffraction Study of the N<sub>2</sub>-CsC<sub>24</sub> System, Molecular Crystals and Liquid Crystals Science and Technology. Section A. Molecular Crystals and Liquid Crystals, 310:1, 99-104

To link to this article: <http://dx.doi.org/10.1080/10587259808045321>

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## Neutron diffraction study of the $N_2$ - $CsC_{24}$ system

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**Abstract :** *Physisorption of molecular  $N_2$  in the 2nd stage  $CsC_{24}$  was studied as a function of temperature and pressure using neutron diffraction. The new structures resulting from reaction were determined from the analysis of the 00l lines and they were related to the intercalation rate of  $N_2$ .*

**Keywords :** graphite, intercalation,  $CsC_{24}$ , Nitrogen, physisorption, neutron diffraction, staging effect.

### INTRODUCTION

Non polar molecules such as nitrogen are known to be physisorbed into the 2nd stage  $MC_{24}$  intercalation compounds at low temperature ( $M=K, Rb$ , and  $Cs$ ) [1]. As reported by Moreh [2], the physisorption of  $N_2$  in  $KC_{24}$  and  $RbC_{24}$  induces the formation of new 2nd stage structures (in which the intercalates species are  $N_2$  and  $Cs$ ) which reversibly segregate from the stage 2 host structure at low temperature. In order to study the new structures resulting from the physisorption of molecular  $N_2$  into  $CsC_{24}$  we performed in situ powder neutron diffraction as a function of pressure and temperature.

### EXPERIMENTAL

$CsC_8$  powder was first prepared by the conventional two bulb method with graphite at  $T_G=250^\circ C$ , and cesium at  $T_{Cs}=200^\circ C$ . A stoichiometric amount of graphite powder (Ceylan, 40 to 60  $\mu m$ ) was added to  $CsC_8$  and the mixture was annealed at  $370^\circ C$  during one week so as to obtain the exact composition  $CsC_{24}$ .

The  $CsC_{24}$  powder was transferred under helium atmosphere into a vanadium container and fixed on a hollow stick. The hollow stick was separated by a valve from a calibrated reservoir connected to a pressure gauge [0-2000 mbar] and a high purity nitrogen bottle allowing the direct measurement of the pressure and the supply with  $N_2$ . The stick was placed in a cryostat working in the range of temperature [300K-10K].

Neutron diffraction patterns were "in situ" recorded at a rate of one data set/five minutes on the D1B diffractometer at  $\lambda=2.52$  Å (ILL, high flux reactor in Grenoble) as a function of pressure and temperature.

Pristine  $\text{CsC}_{24}$  was firstly analyzed under vacuum from 300 K to 10 K and as expected, with neutron diffraction method, no change in the  $00l$  diffraction lines intensities was observed confirming the purity of the sample.

In a conventional experiment, the initial pressure was set to almost 2 bar at room temperature, and the temperature was progressively lowered to 10 K while the pressure was measured at equilibrium at each temperature. The rate of nitrogen intercalation could be estimated from the pressure measurements using the relation  $PV=nRT$  (where  $T$  is the room temperature) with a good approximation at temperature higher than 77 K [3].

In a last experiment,  $\text{CsC}_{24}$  was saturated with liquid nitrogen at a pressure  $>2.7$  bar and at 70 K. The structural evolution of the sample on high pressure of nitrogen was studied in the temperature range [10 K-300 K].

Some of the  $00l$  diffraction lines were fitted using ab-initio calculations and adjusting parameters such as the rate of  $\text{N}_2$  and Cs in the different structures, the atomic Debye Waller factors (C, Cs,  $\text{N}_2$ ), the coherence length along the  $c$  axis, and the proportion and probabilities of occurrence of each layered structures, following a procedure previously used [4], adapted for neutron diffraction.

## RESULT AND DISCUSSION

The intensities of the  $00l$  lines at room temperature before sorption could not be reproduced without taking into account the partial orientation of the GIC crystallites (Fig. 1).

### *$\text{N}_2$ Intercalation at low pressure ( $P \leq 2$ bar)*

When the system was submitted to a pressure close to 2 bar, intercalation started at room temperature. As the temperature decreased, the rate of  $\text{N}_2$  intercalation increased continuously and finally reached the value of 1.5 mole of  $\text{N}_2/\text{CsC}_{24}$  at 100 K (Fig. 2), which is slightly higher than  $\text{N}_2/\text{CsC}_{24}=1.3$  at saturation found by Watanabe et al [1] with an initial pressure of 1 bar.

The intercalation occurs at first randomly in the 2nd stage cesium layers resulting in a great change of  $00l$  lines and a continuous shift of the identity period from 9.35 Å to the maximum value of 9.58 Å at 100 K (Fig. 3).

Below 220 K, a pure first stage  $\text{CsC}_8$  together with the 2nd stage  $\text{Cs}(\text{N}_2)_x\text{C}_n$  ( $n>24$ ) derivative is formed (Fig 4). Due to the intercalation of molecular  $\text{N}_2$ , the internal pressure increased and provoked the motion of a part of cesium in more dense  $\text{CsC}_8$  binary domains [5].

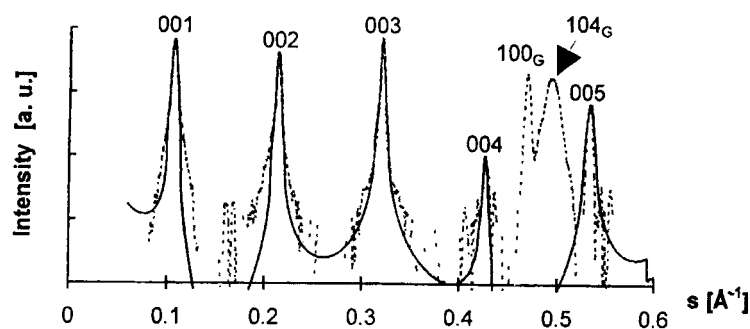


Figure 1 : Logarithmic intensity plot of the experimental scan of pure  $CsC_{24}$  at room temperature (dotted line). The continuous line is the result of a model with a partial preferred orientation function [6] :  $G_1 + (1-G_1)\exp(-G_2 \cdot (\theta-\theta_0)^2)$  [ $G_1=0.54$ ,  $G_2=0.012$ , and  $\theta_0=26^\circ$ ].  $10l$  lines belonging to the graphite network are indexed ( $100_G$  and  $104_G$ ).

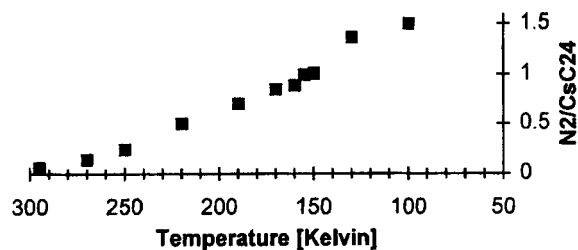


Figure 2 : Ratio of the sorbed molecules at equilibrium to the number of intercalated alkali atoms [ $N_2/CsC_{24}$ ] as a function of the temperature.

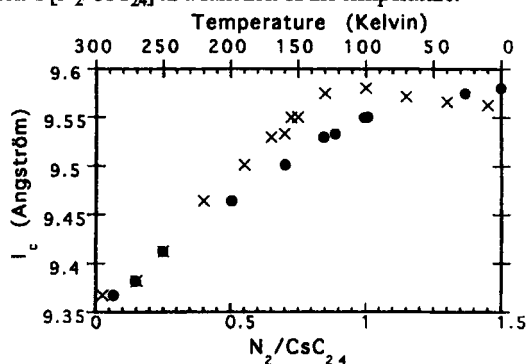


Figure 3 : Variation of the identity period  $I_c$  of the 2nd stage  $Cs(N_2)_x C_n$  ( $n>24$ ) with the amount of sorbed  $N_2$  [ $N_2/CsC_{24}$ ] (circle); with the temperature [in Kelvin] (cross).

The  $00l$  lines of the 2nd stage  $\text{Cs}(\text{N}_2)_x\text{C}_n$  ( $n > 24$ ) were well reproduced at 220 K (Fig. 4) taking  $\text{N}_2/\text{CsC}_{24} = 0.5$  and an  $I_c = 9.51 \text{ \AA}$  in agreement with the experimental values of sorbed amount of nitrogen at this temperature (Fig. 2). The proportion of first stage  $\text{CsC}_8$  is very weak as compared to the 2nd stage  $\text{Cs}(\text{N}_2)_{0.53}\text{C}_{25}$  (Fig. 4).

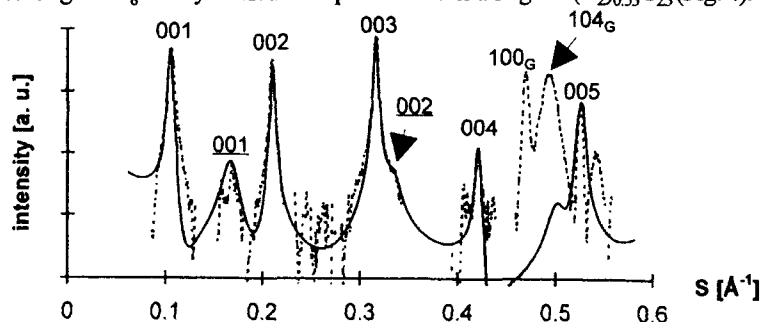


Figure 4 : Logarithmic intensity plot of the experimental scan of  $\text{CsC}_{24}$  saturated with  $\text{N}_2$  ( $\text{N}_2/\text{CsC}_{24} = 0.5$ ) at 220 K (dotted line). The continuous line is the result of a model of the  $00l$  lines with 93% of 2nd stage  $\text{Cs}(\text{N}_2)_{0.53}\text{C}_{25}$  partly interstratified with pure first stage  $\text{CsC}_8$  (underlined  $00l$  indexes).

The identity period  $I_c$  of the 2nd stage  $\text{Cs}(\text{N}_2)_x\text{C}_n$  ( $n > 24$ ) increases linearly with the  $\text{N}_2$  content until 170 K corresponding to a rate of  $0.8 \text{ N}_2/\text{CsC}_{24}$  (Fig 3). At  $T < 170 \text{ K}$ , there is a change of the slope (Fig. 3) and  $I_c$  of the 2nd stage increases more slowly with increasing  $\text{N}_2/\text{CsC}_{24}$  content. Actually below 170 K, the 2nd stage  $\text{N}_2$ -Cs layers are progressively saturated by additional  $\text{N}_2$  molecules so that a first stage  $\text{N}_2$ -Cs ternary graphitide is formed. An example of neutron diffraction pattern at 150 K is shown in figure 5. The 001 peak of this new structure interstratified with the previous  $\text{CsC}_8$  and  $\text{Cs}(\text{N}_2)_x\text{C}_n$  phases is visible but the 002 line is included in the 003 line of the 2nd stage phase, explaining that this peak is larger than the others of the same family (Fig 5). According to a simple model in which the  $\text{N}_2$  molecules would be parallel to the graphite planes, the identity period  $I_c$  of the first stage  $\text{N}_2$ -Cs ternary can be estimated to be  $2r_{\text{N}_2} + d_G = 6.45 \text{ \AA}$  (where  $r_{\text{N}_2}$  and  $d_G$  are respectively the Van der Waals radii of Nitrogen :  $1.55 \text{ \AA}$  and the distance between two graphite layers :  $3.35 \text{ \AA}$ ) in good agreement with the experimental value of  $6.4 \text{ \AA}$ . However, the measured value of  $2r_{\text{N}_2}$  was found to be  $2.86 \text{ \AA}$  [2] yielding  $2r_{\text{N}_2} + d_G = 6.21 \text{ \AA}$  so that the measured identity period ( $6.4 \text{ \AA}$ ) could be explained by the tilt of the nitrogen molecules in relation to the graphite planes.

At 100 K, complete physisorption is achieved. Indeed, from  $T = 100 \text{ K}$  to  $10 \text{ K}$  the  $00l$  intensities remain unchanged, and only the identity period of the 2nd stage slowly decreases with decreasing temperature (Fig 3). Due to the freezing of the cesium atoms at low temperature [4] the motion of the cesium atoms cannot occur

anymore, and as a consequence the system remains a mixture of the three structures previously described. However changes were observed in the  $hk$  lines (Fig. 5) appearing at  $T < 160$  K, attributed to the 2D order of nitrogen and cesium.

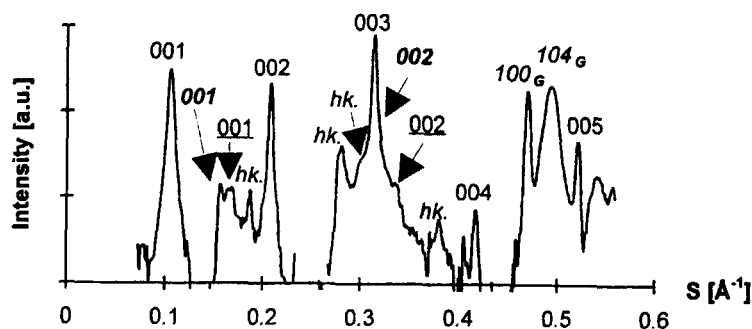


Figure 5 : Experimental scan of  $CsC_{24}$  saturated with  $N_2$  at 150 K. The pattern consists of  $00l$  lines of a first stage  $CsC_8$  (underlined indexes) and a first stage  $Cs(N_2)_x C_m$  (bold indexes) with the lines belonging to the 2nd stage  $Cs(N_2)_x C_n$  (normal indexes).

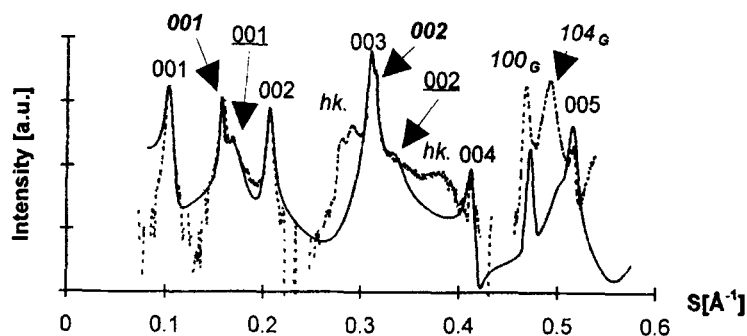


Figure 6 : Logarithmic intensity plot of an experimental scan of  $CsC_{24}$  saturated with  $N_2$  at 160 K (dotted line). The continuous line is the result of a model of the  $00l$  lines with a mixture of 25% first stage  $Cs(N_2)_{1.8}C_{16}$  ( $l_c=6.37$  Å, bold indexes) and 17% of the binary first stage  $CsC_8$  ( $l_c=5.98$  Å, underlined indexes) interstratified in the 2nd stage  $Cs(N_2)_{2.1}C_{32}$  ( $l_c=9.71$  Å, normal indexes).

#### Intercalation at higher pressure

When high pressure ( $>2.7$  bar) was applied on the system cooled at 70 K, part of the nitrogen was physisorbed in the 2nd stage layers and liquid nitrogen was also formed. When lowering temperature, reversible liquid-solid transition of

nitrogen was observed. Upon heating, the nitrogen pressure applied to the system increase dramatically leading above  $T=150$  K to the formation of a first stage  $N_2$ -Cs ternary graphite ( $I_c=6.37$  Å) together with 2nd stage  $Cs(N_2)_x C_n$  ( $n>24$ ) and pure stage one  $CsC_8$ . When the system is submitted to more than 9 bar at  $T=165$  K, the content of first stage derivative was found about seven times higher than at 1.3 bar. This result indicates that the proportion of this first stage  $N_2$ -Cs ternary graphite is related to the pressure and as a consequence, the rate of  $N_2$  intercalated increases with increasing pressure applied to the system. Fitting the  $00l$  lines intensities (Fig. 6), it can be deduced that the compound at 160 K is a mixture of 25% segregated first stage  $Cs(N_2)_{1.8}C_{16}$  ( $I_c=6.37$  Å) and 17% binary first stage  $CsC_8$  ( $I_c=5.98$  Å) interstratified with the 2nd stage  $Cs(N_2)_{2.1}C_{32}$  ( $I_c=9.71$  Å). The first stage ternary graphite structure reversibly disappears at 170 K and only 2nd stage  $Cs(N_2)_x C_n$  ( $n>24$ ) and pure  $CsC_8$  remain stable above this temperature.

### CONCLUSION

In agreement with previous results, it was found that due to its larger interlayer distance,  $CsC_{24}$  is a best sorptive agent than  $RbC_{24}$  and  $KC_{24}$  [1-2]. As decreasing the temperature, the  $N_2$  molecules firstly intercalates in the 2nd stage  $Cs(N_2)_x C_n$  ( $n>24$ ) and at  $T<170$  K a first stage ternary  $Cs-N_2$  graphite is formed. The sorption phenomena is enhanced by applying high pressure and it is completely achieved at  $T=100$  K. Upon decreasing the temperatures, the  $N_2$  molecules firstly intercalates in the cesium layers. Pressure effect leads to the redistribution of cesium giving a mixture of 2nd stage  $Cs(N_2)_x C_n$  ( $n>24$ ) and 1st stage  $CsC_8$ . Below 170 K, increasing the amount of intercalated nitrogen leads to transferring a part of the 2nd stage  $Cs(N_2)_x C_n$  into 1st stage  $Cs(N_2)_y C_m$ . At 100 K, the mobility of the cesium atoms becomes so low that further intercalation cannot proceed. Applying high nitrogen pressure enhances drastically the amount of 1st stage ternary graphite which is obviously the ultimate step of nitrogen intercalation.

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